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Investigation of Bulk Traps by Conductance Method in the Deep Depletion Region of the Al₂O₃/GaN MOS Device

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Abstract

Conductance method was employed to study the physics of traps (e.g., interface and bulk traps) in the Al₂O₃/GaN MOS devices. By featuring only one single peak in the parallel conductance (G_p/ω) characteristics in the deep depletion region, one single-level bulk trap ($E_C=0.53$ eV) uniformly distributed in GaN buffer was identified. While in the subthreshold region, the interface traps with continuous energy of $E_C=0.4\sim0.57$ eV and density of $0.6\sim1.6 \times 10^{12} \text{ cm}^{-2}$ were extracted from the commonly observed multiple G_p/ω peaks. This methodology can be used to investigate the traps in GaN buffer and facilitates making the distinction between bulk and interface traps.

Keywords: Al₂O₃/GaN MOS channel device, Conductance method, Buffer traps, Interface traps

Background

Owing to the superior properties of high electron mobility, high breakdown voltage, high-power density, low on-resistance, and high temperature operation capability, GaN heterojunction field-effect transistors (HFETs) have been considered as a promising solution for next-generation energy-efficient power electronics and attracted tremendous attention in the last two decades [1]. For power switching applications, the enhancement-mode (E-mode) transistors are highly preferred rather than the depletion-mode (D-mode) devices for the inherent fail-safe operation and simple gate driver circuitry. Despite the various technologies proposed to realize E-mode, GaN HFETs such as p-cap gate [2, 3], fluorine plasma ion implantation [4], and cascode technology [5], the MOSFET with partially or fully recessed gate is considered as a promising candidate because of its high-threshold voltage (V_{TH}), large gate swing for improved fail-safe capability [6, 7], and low on-resistance [8]. Moreover, the MOS-gate is compatible with the mainstream gate driver ICs. However, the traps (e.g., interface and bulk traps) tarnish the advantages of GaN HFETs due to the stability and reliability issues such as

V_{TH} instability [9], drain lag or gate lag [10], and power slump. Besides the surface/interface traps, the GaN power HFETs' performance such as the breakdown voltage and dynamic on-resistance could be substantially affected by the bulk traps in GaN buffer layer in high-voltage-switching applications [11, 12] since the high electric-field is prone to trigger the buffer traps for dynamic charging/discharging. Hence, it is of great significance to characterize the buffer traps of GaN MOS devices.

Bulk traps in GaN MOS devices have been studied by deep-level transient spectroscopy (DLTS) [13] and pulse measurement [14]. However, though the dynamic charge/discharge process of both bulk and interface trap-induced transient behavior may simultaneously appear in the same spectrum, extra effort is required to differentiate between the bulk and interface traps when using DLTS-like techniques and pulse measurement [15, 16]. On the other hand, the conductance method has been widely used to evaluate the interface traps in AlGaN/GaN MIS structures as well as GaN-based MIS structures [17–19]. Moreover, it is possible to discriminate the bulk and interface traps in the conductance method by studying its bias dependence because the conductance loss is sensitive to the traps within a few kT/q around the Fermi level. In this letter, the conventional conductance method normally used to characterize the interface traps is employed to study the bulk traps (BT) in GaN buffer for the first time. Two trap-

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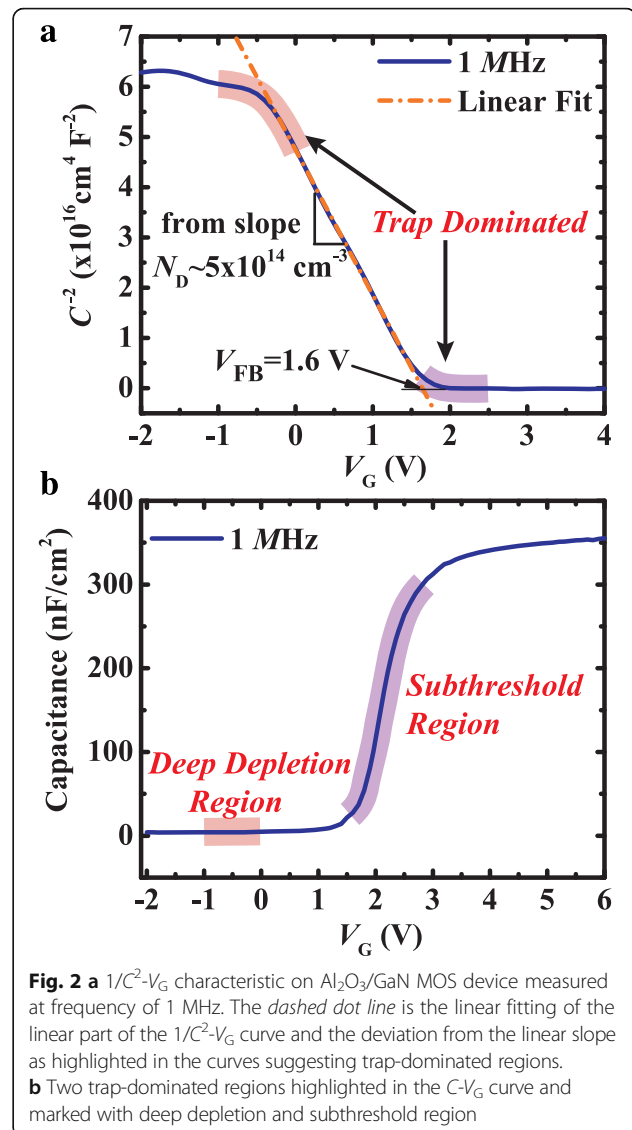
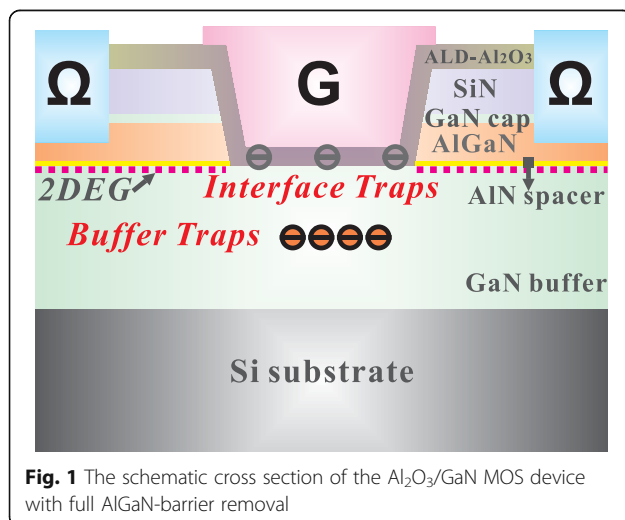
dominated regions were found in the $\text{Al}_2\text{O}_3/\text{GaN}$ MOS structure with full barrier recess. In the deep depletion region, only one single Gp/ω peak is captured at the measured bias voltage ranging from -1 to 0 V, revealing the bulk traps with a single level in GaN buffer. The energy of the bulk trap was determined to be $E_C-0.53$ eV. While in the subthreshold region, the interface traps with continuous energy levels that result in multiple Gp/ω peaks were observed within the measured bias range of $1.6\sim 2.6$ V. The energy levels were extracted to be in the range of 0.4 to 0.57 eV below GaN conduction band (CB).

Methods

The device structure used in this work is shown in Fig. 1. The $\text{Al}_2\text{O}_3/\text{GaN}$ MOS device is fabricated on a commercial $\text{Al}_{0.25}\text{Ga}_{0.75}\text{N}/\text{GaN}$ heterostructure grown on a 4-in Si (111) substrate by MOCVD. The AlGaN layer was fully recessed by using the low-damage hybrid recess technique. The detail fabrication process can be found in our previous work [7]. A 20-nm Al_2O_3 layer was deposited by atomic layer deposition (ALD) as the gate oxide. The conductance-frequency (G - f) and capacitance-voltage (C - V) characteristics were measured by an Agilent B1500A Semiconductor Device Analyzer equipped with a Cascade probe station. The frequency-dependent conductance measurements are shown within the range of 1 to 5 MHz.

Results and Discussion

The $1/C^2$ - V_G and C - V_G characteristics of the $\text{Al}_2\text{O}_3/\text{GaN}$ MOS device measured at frequency of 1 MHz are shown in Fig. 2. A linear region was observed in the $1/C^2$ - V_G curve in the gate bias voltage range of $0\sim 1.5$ V in Fig. 2a. The flat band voltage (V_{FB}) is determined to be 1.6 V by linear extrapolation of the $1/C^2$ - V_G curves to the intercept with the abscissa [20]. The background-doping concentration of GaN buffer N_D was extracted to be $5 \times$



10^{14} cm^{-3} from the linear slope. The good linear fitting suggests negligible trap states in the depletion region, since the linearity of the slope is strongly affected by the charge/discharge process of the traps [21, 22]. However, a decrease in slope implies the existence of residual trap states with the applied gate voltage both lower than 0 V and higher than 1.5 V, which corresponds to the deep depletion region and subthreshold region as highlighted in Fig. 2b.

In order to further differentiate the types of the traps (e.g., interface trap or bulk trap) and study the trap characteristics including the trap levels and densities, f -dependent conductance (Gp/ω) measurements were performed. The Gp/ω as a function of radial frequency ($\omega = 2\pi f$) can be correlated to the traps density D_T and trap response time τ_T . For the case of bulk trap states with discrete energy level, Gp/ω is given by [23, 24]

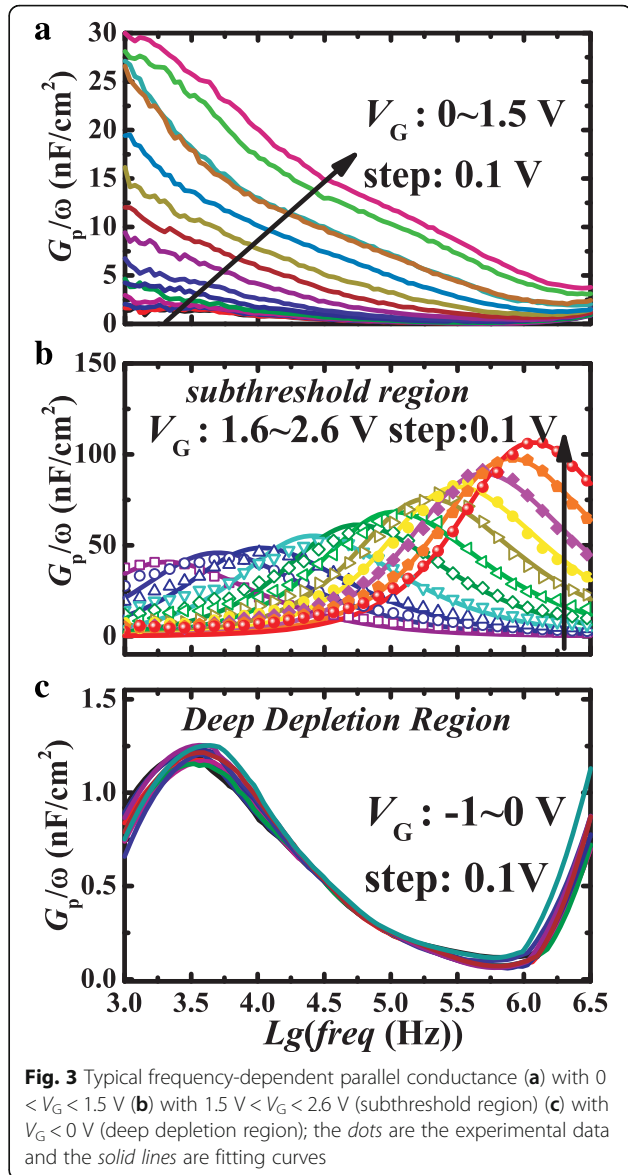
$$\frac{G_p}{\omega} = \frac{q\omega\tau_T D_T}{1 + \omega^2\tau_T^2}, \quad (1)$$

whereas, for interface trap states with distributed energy levels, G_p/ω is given by

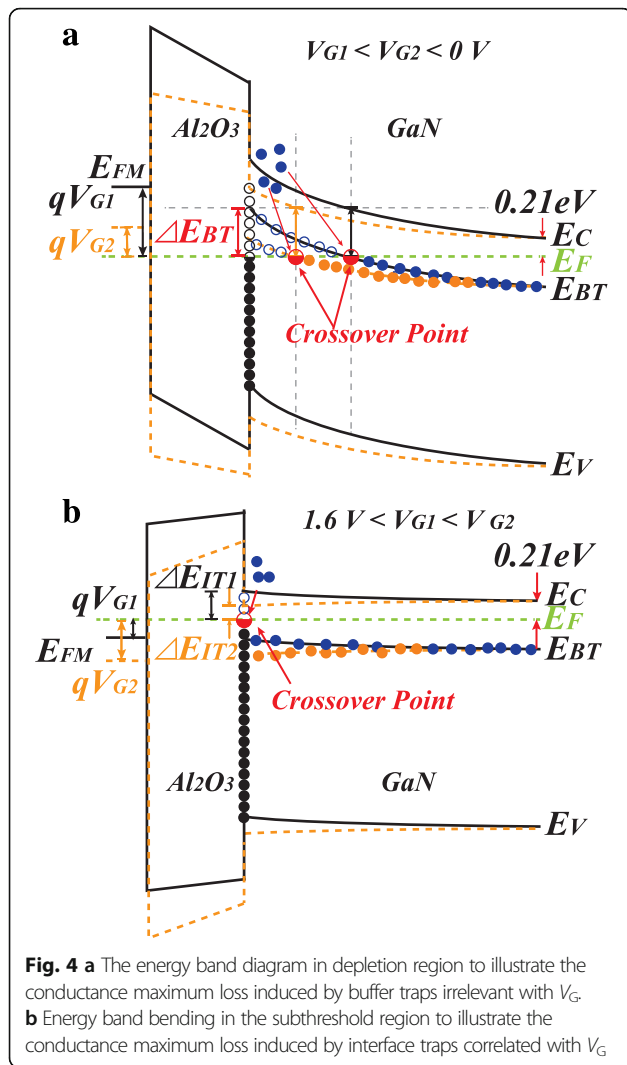
$$\frac{G_p}{\omega} = \frac{qD_T}{2\omega\tau_T} \ln(1 + \omega^2\tau_T^2). \quad (2)$$

The maximum loss can be obtained when the trap states are in resonance with the applied AC signal for dynamic discharging/charging, which occurs when the trap states are exactly half-filled, i.e., when the energy level of the trap states crosses with the semiconductor Fermi level [24]. Consequently, the trap time constant corresponds to the maximum of G_p/ω and can be determined by setting the derivative $\partial(G_p/\omega)/\partial(\omega\tau_T)$ to zero. By following the above two equations, the $\omega\tau_T$ is found to be 1 and 1.98 for bulk traps and interface traps, respectively. Thus, the G_p/ω peak frequency is associated with the trap energy level and the peak value is related to the trap density. The G_p/ω curves monotonically decrease with the increasing frequency without G_p/ω peaks at $0 \text{ V} < V_G < 1.5 \text{ V}$ (see Fig. 3a) that corresponds to the observed linear regime shown in Fig. 2a, which reinforces the negligible charge/discharge processes of traps in this region. On the other hand, as shown in Fig. 3b, multiple peaks were observed in the G_p/ω curves while the device operated in the subthreshold region as the applied bias above the flat band voltage. In this region, the peaks steadily shift to higher frequencies with the increasing applied voltage, which is the typical G_p/ω characteristic indicating the presence of interface traps with continuous energy levels as commonly observed in conventional conductance measurements [17–19, 23]. The G_p/ω characteristics in the deep depletion region ($-1 \text{ V} < V_G < 0 \text{ V}$) are plotted in Fig. 3c which exhibits quite a difference compared with that observed in the subthreshold region. The G_p/ω curves within the measured bias range featured an identical profile with a single peak and the same G_p/ω peak value, which suggests the existence of trap state with only one single energy level while the electron emission rate of the trap state is irrelevant with V_G . It is well known that bulk traps stem from the defects, and impurities are usually in a uniform-distribution throughout the bulky semiconductor. Correspondingly, the bulk traps are in discrete energy levels and capable of inducing the same loss peaks at the same frequency under various biases [24]. Hence, the G_p/ω characteristic with one peak observed in the deep depletion region reinforced that the bulk trap with one single level in GaN buffer layer was identified by conventional conductance measurement.

The two types of traps (e.g., bulk and interface traps) can be easily discriminated by studying the bias-



dependent conductance in different operation regions of the device as illustrated in the energy bandgap diagram of $\text{Al}_2\text{O}_3/\text{GaN}$ MOS device in Fig. 4. The conductance method is sensitive to traps within a few kT/q around the Fermi level marked as the crossover point in Fig. 4. As shown in Fig. 4a, the single-level bulk traps are uniformly distributed in GaN buffer. The magnitude of G_p/ω peak originates from the bulk traps not varying with gate bias regardless of its spatial location. Meanwhile, the interface traps that cross over the Fermi level are much deeper in the GaN bandgap than the bulk trap at the given bias voltage. Therefore, the conductance loss was dominated by the dynamic response of the bulk traps rather than the interface traps in the deep depletion region. Further sweep up the bias voltage (e.g., $0 \text{ V} < V_G < 1.5 \text{ V}$), all of the bulk traps were filled with



electrons as shown in Fig. 4b. On the other hand, the interface traps with a wide energy distribution may still capture the electrons. However, the conductance loss was not detected at the measured frequencies (e.g., 1 kHz–5 MHz) due to the extremely large time constant associated with the deep levels of the interface traps. Consequently, the $1/C^2$ - V_G plot exhibits a linear characteristic and the Gp/ω curves show a monotone property while $0 < V_G < 1.5$ V as shown in Fig. 2a and Fig. 3a, respectively. In the subthreshold region (1.6 V $< V_G < 2.6$ V), it can be seen from Fig. 4b that the conductance loss is solely contributed by the charging/discharging of the interface traps. Because the energy level of the interface traps that cross over the Fermi level are shallower in the subthreshold region, the relatively small time constant for dynamic charging/discharging can be detected as the multiple Gp/ω peaks measured in Fig. 3b.

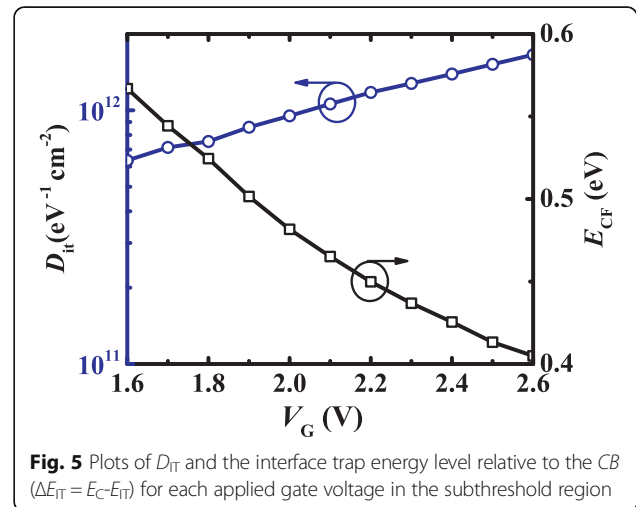
By using Eq. (2), the fitting curves exhibit good agreement with the measured Gp/ω characteristics in the

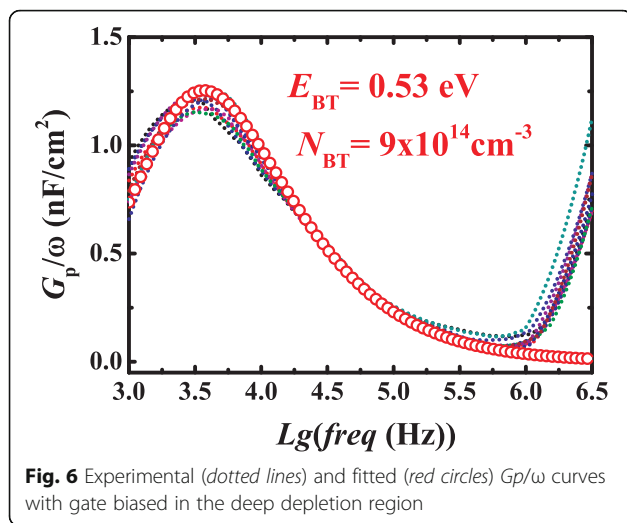
subthreshold region. The interface trap density and energy distribution profile were extracted and plotted versus the applied voltage in Fig. 5. The trap energy levels relative to the CB ΔE_T as function of response time τ_T given by the Shockley–Read–Hall statistics were extracted using the equation [14]:

$$\Delta E_T = k_B T \ln(v_{th} \sigma_n N_C \tau_T), \quad (3)$$

where the capture cross section of the trap $\sigma_n = 4 \times 10^{-13}$ cm², the electron thermal velocity $v_{th} = 2.6 \times 10^7$ cm/s, the density of states at GaN CB $N_C = 2.2 \times 10^{18}$ cm⁻³, the Boltzmann constant $k_B = 1.38 \times 10^{-23}$ J/K, and temperature $T = 300$ K were used [25]. The interface trap levels are in the range of 0.4 to 0.57 eV below GaN CB with D_{IT} decreased from 1.6×10^{12} to 0.6×10^{12} cm⁻².

Similarly, the energy level and density of bu traps in GaN buffer also can be extracted by fitting the Gp/ω characteristics with Eqs. (1) and (3) as shown in Fig. 6. More importantly, as the measured Gp/ω characteristics exhibit an identical profile with only one single peak value at various bias voltages, the measured data were well fitted by a single curve instead of a series Gp/ω curves for the interface trap that feature a continuum energy-level distribution. Thus, single-level bulk traps with sheet density 1.5×10^{10} cm⁻² and $E_{BT} = E_C - 0.53$ eV (close to the reported Fe-induced level at 0.5 ± 0.1 eV below the GaN conduction band edge in GaN buffer [26, 27]) was extracted. Accordingly, the volume density N_{BT} of the buffer trap can be obtained by dividing the sheet density by Debye length $L_D = \sqrt{kT\epsilon_S\epsilon_0/(q^2N_D)}$ as the depletion depth being in the same order of magnitude or smaller than the Debye length [28]. With the unintentional doping density $N_D = 5 \times 10^{14}$ cm⁻³ extracted from the slope of the $1/C^2$ - V_G curve, N_{BT} was estimated to be 9×10^{14} cm⁻³.





Conclusions

In conclusion, for the first time, the conventional conductance method was used to study the buffer traps in the $\text{Al}_2\text{O}_3/\text{GaN}$ MOS device with full barrier removal. The bulk traps with a single energy level and uniformly distributed in GaN buffer that leads to a single G_p/ω peak were detected by f -dependent conductance measurements in the deep depletion region. On the other hand, the interface traps with wide energies were measured in the subthreshold region, which corresponds to the multiple G_p/ω peaks observed in the f -dependent conductance measurements. Due to the different f -dependent conductance response originating from the different energy and spatial distributions, the demonstrated approach is much easier to be used to investigate the physics of the bulk traps in GaN buffer.

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Authors' contributions

YYs jointly conceived the study with ZJL. YYS, ABZ, LYZ, and YS performed the experiments. YYS and ZJL performed all the data analyses and wrote the original draft of the manuscript. QZ, WJC, and BZ reviewed and edited the manuscript. All authors reviewed the manuscript. All authors read and approved the final manuscript.

Competing interests

The authors declare that they have no competing interests.

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